

# Refined 0D–3D dynamic cluster model of magnetic susceptibility in $\text{Ge}_{1-x-y}\text{Sn}_x\text{Mn}_y\text{Te}$ : the role of secondary phase microinclusions

V. E. Slynko<sup>1</sup>, M.V. Tovarnitskii<sup>1</sup>, A.V. Zasloukin<sup>1</sup>, V.V. Netyaga<sup>1</sup>, L. Kilanski<sup>2</sup>, S. Piotrowska<sup>2</sup>

<sup>1</sup>Chernivtsi Branch of Frantsevych Institute for Problems of Materials Science, NASU,  
5 Iryny Vilde St., Chernivtsi, 58001, Ukraine

<sup>2</sup>Institute of Physics, Polish Academy of Sciences, Aleja Lotnikow 32/46, PL-02668 Warsaw, Poland  
e-mail: slynko.vasily@gmail.com

The complex nature of the magnetic response in diluted magnetic semiconductors (DMS)  $\text{Ge}_{1-x-y}\text{Sn}_x\text{Mn}_y\text{Te}$ , particularly the anomalous double-peaked temperature dependence of the dynamic magnetic susceptibility  $\chi_{ac}(T)$ , has been previously described within the framework of a dynamic 0D–3D cluster magnetic subsystem model [1]. While the initial theoretical approach successfully parameterized the macroscopic magnetic behavior, the microscopic mechanism driving the inter-cluster interactions requires refinement to avoid overestimating the role of bulk thermal lattice vibrations.

Recent comprehensive micro-XRF mapping of DMS  $\text{Ge}_{1-x-y}\text{Sn}_x\text{Mn}_y\text{Te}$  crystals (specifically, samples from ingot 1292, Figs. 1-2) has revealed the physical origin of these magnetic anomalies. The data unambiguously demonstrate the presence of structural micro-inhomogeneities. Instead of an idealized homogeneous matrix, the material contains localized microregions with pronounced deviations from stoichiometry, primarily characterized by an enriched tellurium content (reaching 65.0–75.0 at.% Te, in stark contrast to the nominal ~50 at.% of the host lattice). These secondary phase microinclusions exhibit typical dimensions ranging from 5 to 15  $\mu\text{m}$  and occupy a volumetric fraction of approximately 5–7%.

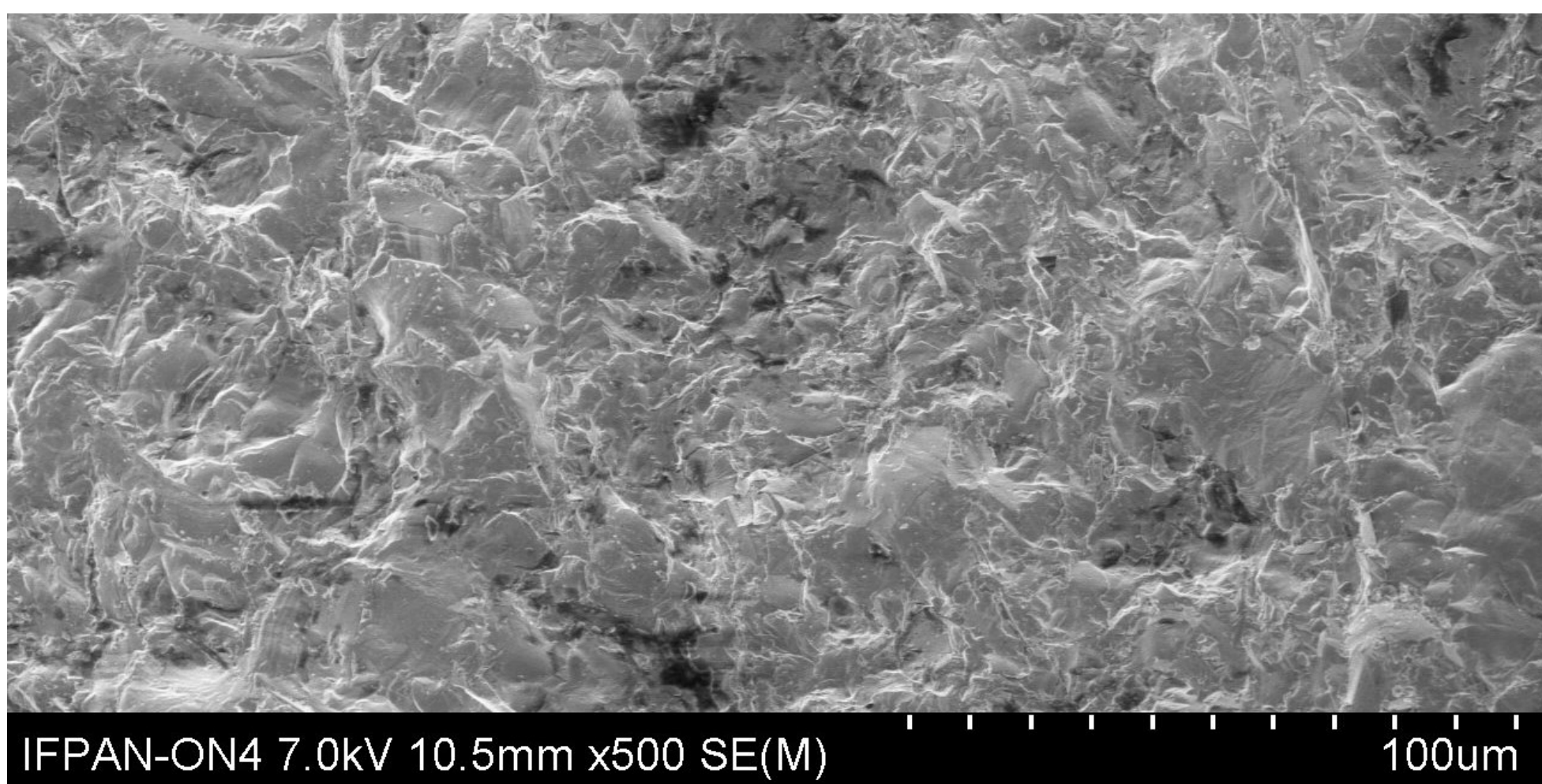


Fig. 1. Secondary phase microinclusions (enriched Te) in the 1292/2  $\text{Ge}_{1-x-y}\text{Sn}_x\text{Mn}_y\text{Te}$  ( $\times 500$ ).

The formation of these Te-enriched microinclusions is an inevitable consequence of the crystal growth thermodynamics. Specifically, they arise due to compositional instabilities and segregation processes in the melt [2], which lead to microscopic fluctuations in the composition and density of *low-temperature eutectics* throughout the entire crystallization process. Therefore, absolute spatial

homogeneity in these quaternary solid solutions is thermodynamically unattainable.

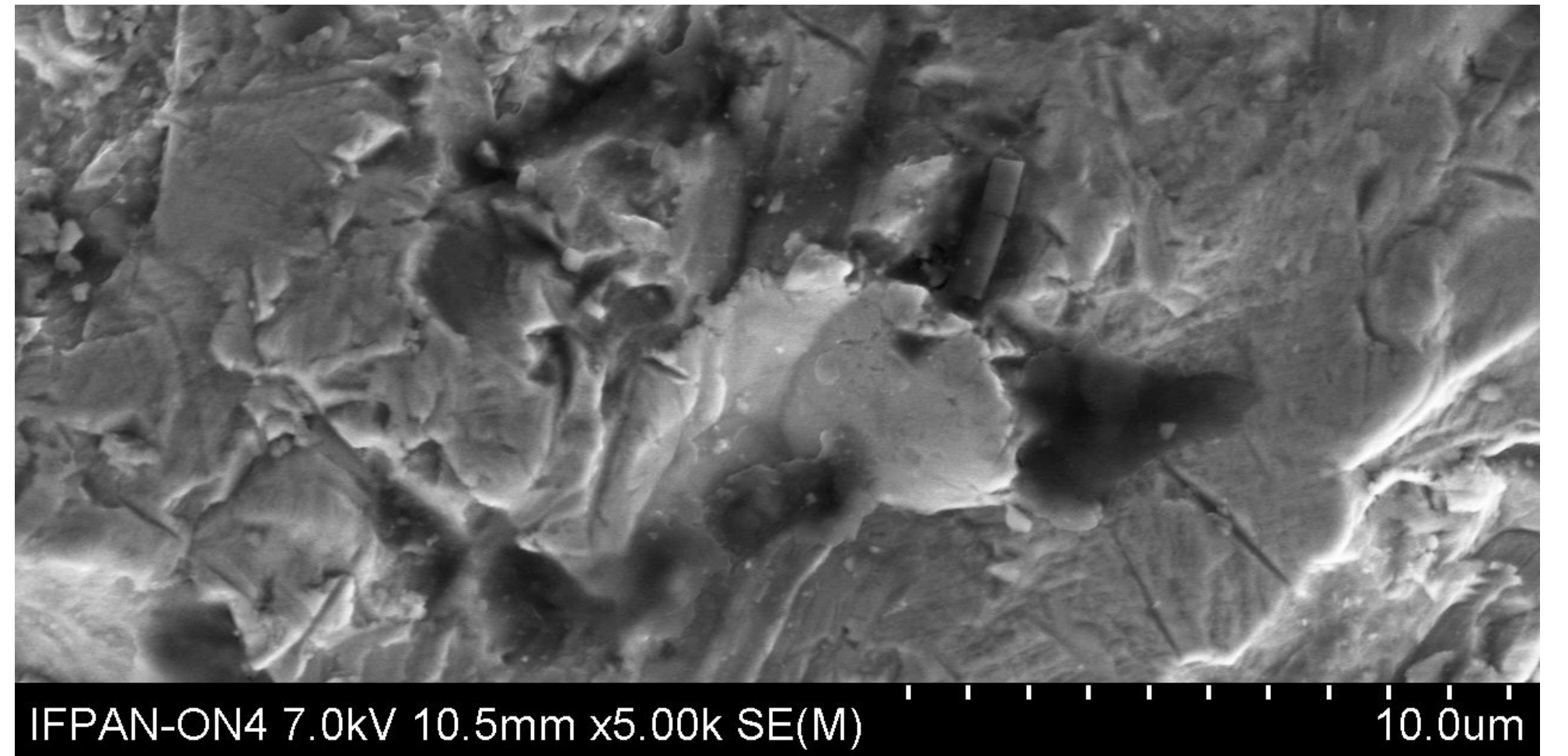


Fig. 2. Secondary phase microinclusions (enriched Te) in the 1292/2  $\text{Ge}_{1-x-y}\text{Sn}_x\text{Mn}_y\text{Te}$  ( $\times 5000$ ).

Recognizing this intrinsic micro-heterogeneity allows for a critical enhancement of the 0D–3D cluster model [1]. The primary driver of the complex magnetic response  $\chi_{ac}(T)$  is not the uniform thermal vibration of an ideal lattice, but rather the interface effects at the boundaries between the Te-enriched microinclusions and the main matrix. These boundaries act as intense centers for phonon scattering. Furthermore, the local lattice strain at these interfaces causes a sharp, exponential dependence of the p-d exchange interaction constant ( $J_{pd}$ ) on the overlap of wave functions [3]. Even minute displacements of magnetic ions ( $\text{Mn}^{2+}$ ) within these highly strained boundary regions result in significant modulations of the local magnetic order.

Thus, the refined 0D–3D cluster model, grounded in direct micro-XRF evidence, demonstrates that the macroscopic magnetic anomalies in  $\text{Ge}_{1-x-y}\text{Sn}_x\text{Mn}_y\text{Te}$  are governed by the dynamic interplay and boundary scattering effects between structurally distinct, thermodynamically inevitable micro-phases, rather than purely by uniform thermal lattice dynamics.

[1] V. Slynko et al., "Dynamic cluster magnetic subsystems in diluted magnetic semiconductor  $\text{Ge}_{1-x-y}\text{Sn}_x\text{Mn}_y\text{Te}$ ", *Low Temp. Phys.* 52, 262 (2026). DOI: 10.1063/10.0042378

[2] P. Rudolph, M. Mühlberg, "Basic problems of vertical Bridgman growth of CdTe", *Mater. Sci. Eng. B* 16 (1-3), 8-16 (1993). DOI: 10.1016/0921-5107(93)90005-8

[3] A. Abragam, B. Bleaney, *Electron Paramagnetic Resonance of Transition Ions*, Oxford Univ. Press (1970). ISBN: 978-0198512509.