Hydrostatic pressure effect on the pseudogap in slightly doped $Y_{0.77}Pr_{0.23}Ba_2Cu_3O_{7-\delta}$ single crystals E. V. Petrenko¹, L. V. Bludova¹, A. S. Kolesnik¹, A. Sedda², E. Lähderanta², R. V. Vovk³, A. L. Solovjov^{1,2,4} ¹B.Verkin Institute for Low Temperature Physics and Engineering of NAS of Ukraine, 47 Nauky Ave., Kharkiv, 61103, Ukraine ²Lappeenranta University of Technology, School of Engineering Science, 53850 Lappeenranta, Finland ³Physics Department, V. N. Karazin Kharkiv National University, Svobody Sq. 4, 61022 Kharkiv, Ukraine ⁴Institute for Low Temperatures and Structure Research, of PAS, Wroclaw 50-422, Poland BludovaLV2102@gmail.com



Introduction and motivation

The discovery of high-temperature superconductors (HTSCs) is undoubtedly one of the landmark events in modern solid state physics. However, despite the efforts of numerous scientific groups and an extraordinary number of publications on HTSCs, the mechanism of superconducting (SC) pairing, which makes it possible to obtain Cooper pairs at T >> 100 K, is still not clear. The pseudogap (PG) state, which is opened in cuprate HTSCs below the characteristic temperature $T^* >> T_c$, is one of the most mysterious and simultaneously interesting phenomena in modern solid state physics. It is well established that in HTSCs, the PG is observed when the charge carrier concentration varies between slightly doped (SD) and optimally doped levels. Understanding the PG physics would definitely shed more light on the mechanism of superconducting pairing in HTSCs, which is also not fully clarified yet. The $YBa_2Cu_3O_{7-\delta}$ cuprate is believed to be the most reliable material for studying the PG, especially when high pressures is applied.

Temperature dependences of the resistivity of slighly doped $Y_{0.77}Pr_{0.23}Ba_2Cu_3O_{7-\delta}$ Dependences of excess conductivity $\sigma'(T)$ versus the reduced temperature ε in double logarithmic scale for S1 and S5 in the range of SC fluctuations near T_c single crystals ($T_c \approx 66.4$ K) at different pressure. The arrows designate T^* . in comparison with fluctuation theories: 350 **3D** Aslamazov-Larkin (AL) **2D Maki-Thompson (MT)** $\rho_N(T) = \alpha T + \beta$ $\alpha, \beta - const$ • S1 300 (red straight lines) (blue curves) T*=163 $\sigma_{AL}' = \frac{e^2}{32\hbar\xi_c} \varepsilon^{-1/2} \qquad \qquad \sigma_{MT}' = \frac{e^2}{8d\hbar} \cdot \frac{1}{1 - \alpha / \delta} \cdot \frac{1}{1 - \alpha / \delta} \cdot \frac{1 + \alpha + \sqrt{1 + 2\alpha}}{1 + \delta + \sqrt{1 + 2\delta}} \cdot \varepsilon^{-1}$ **E** 250 • S5 $\sigma'(T) = \frac{1}{\rho(T)} - \frac{1}{\rho_N(T)}$ G₂₀₀



Table. Resistive and fluctuation conductivity parameters (FLC) parameters of $YPrBa_2Cu_3O_{7-\delta}$ single crystal at different presurre.

| GPa | ρ(300K) | ρ(100K) | T _c | T_c^{mf} | To | <i>T</i> ₀₁ | T _G | <i>d</i> ₀₁ | $\xi_c(0)$ | Gi |
|------|-----------------|-----------------|----------------|--------------|--------------|------------------------|----------------|------------------------|------------|--------|
| | $\mu\Omega(cm)$ | $\mu\Omega(cm)$ | (K) | (K) | (K) | (K) | (K) | (A) | (A) | |
| 0 | 331.24 | 141.16 | 66.4 | 67.55 | 68.0 | 74 | 67.6 | 3.18 | 1.00 | 0.0007 |
| 0.41 | 296.9 | 124.6 | 68.25 | 69.18 | 70.0 | 76.4 | 69.3 | 3.88 | 1.2 | 0.0017 |
| 0.64 | 275.6 | 113.2 | 68.91 | 69.85 | 70.2 | 75.4 | 69.9 | 3.18 | 0.90 | 0.0007 |
| 0.87 | 259.3 | 108.3 | 69.07 | 70.18 | 70.6 | 75.7 | 70.3 | 3.3 | 0.89 | 0.0017 |
| 1.1 | 247.36 | 102.14 | 69.7 | 70.55 | 70.9 | 76.1 | 70.7 | 3.41 | 0.884 | 0.0021 |

L. G. Aslamazov and A. L. Larkin, Phys. Lett., 26A, 238 (1968).





Main results

$$\Delta^*(T) = T \ln \frac{\tau^2 \left(T \right)}{\sigma'(T) 16\hbar \xi_c(0) \sqrt{2\varepsilon^*_{c0} \sinh(2\varepsilon/\varepsilon^*_{c0})}}$$



Conclusions

In our work, for the first time, we carried out the analysis of the influence of hydrostatic pressure up to 1.1 GPa on the temperature dependence of pseudogap $\Delta^*(T)$ of the SD $Y_{0.77}$ Pr_{0.23}Ba₂Cu₃O₇₋₈ single crystals.

It is shown that the pressure effect on T_c and resistivity $\rho(T)$ is different. Under pressure, $\rho(T)$ decreases, while T_c increases, which is associated with the redistribution of charge carriers in the CuO₂ planes.

It is demonstrated that under pressure, the $\Delta^*(T)$ increases with a rate $d\ln\Delta^*/dP = 0.9$ GPa⁻¹, which is most likely due to a decrease in the frequencies of the phonon spectrum of the superconductor when pressure is applied.

It is revealed that without pressure, a "magnetic" maximum occurs on the $\Delta^*(T)$ at high temperatures, followed by a linear section with a positive slope, limited by the temperatures of structural transition T_s and spin density wave ordering T_{SDW} . At 1.1 GPa, the maximum disappears.