

# Intrinsic and pinning anisotropy from surface impedance measurements in superconducting YBCO and Fe(Se,Te) thin films

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The sources of anisotropy in layered type II superconductors  $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$  (YBCO) and Fe-based compound like  $\text{FeSe}_x\text{Te}_{1-x}$  (FeSeTe) have different physical origin: the anisotropic electron effective mass (giving rise to anisotropic coherence length) determines what we call here *intrinsic anisotropy*  $\gamma$ . Extended structures, including the layered structure itself, twin planes, grain boundaries, columnar or some artificial pinning centres (APC), introduce a complex vortex pinning landscape and then they give rise to the *pinning anisotropy*.

Using microwave measurements of the complex surface impedance  $Z=R+iX$  in a dc magnetic field we are able to distinguish the intrinsic from pinning anisotropy. In particular, the changes in  $Z(H)$  are due to vortex drag (the origin of flux-flow resistivity  $\rho_{\text{ff}}$ , which depends on the anisotropy through the effective vortex mass), to the recall of the vortex from the pinning centre (given by the pinning constant  $k_p$ ), and to flux-creep jumps (quantified by a normalized dimensionless parameter  $\chi$ ) [1]. Measurements of  $Z$  at different dc magnetic fields  $H$  and field orientations  $\theta$ , and at different frequencies, allow to extract  $\rho_{\text{ff}}$ ,  $k_p$  and  $\chi$ , whence the intrinsic and pinning anisotropies [2].

We measure  $Z(H,\theta)$  using a dielectric-loaded resonator technique in different setups [3], at 16, 27 and 47 GHz, up to a dc field of 1.3 T. We present results obtained in YBCO thin films grown by Pulsed Laser Deposition (PLD) [4], and in pristine  $\text{FeSe}_x\text{Te}_{1-x}$  thin films grown by PLD from  $\text{FeSe}_{0.5}\text{Te}_{0.5}$  targets [5]. We extract the intrinsic anisotropy  $\gamma$ , obtaining in YBCO consistent values  $\gamma = 5.0 \pm 0.5$  [6], and  $\gamma \approx 2$  in FeSeTe. On the other hand, we obtain very different pinning anisotropies in all compounds. In YBCO, the defects affect in a complex fashion the anisotropic  $k_p$ , while in FeSeTe  $k_p(\theta)$  does not show significant departures from the effective mass effect. Interestingly, we do not find signatures of the multigap nature of FeSeTe in the angular dependence of the flux-floe resistivity: a single-band scaling approach [7] applies.

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